

# Landau gauge Yang-Mills correlation functions

Anton Konrad Cyrol

Ruprecht-Karls-Universität Heidelberg

based on

- AKC, Fister, Mitter, Pawłowski, Strodthoff, PRD, arXiv:1605.01856 [hep-ph]
- AKC, Mitter, Strodthoff, FormTracer, arXiv:1610.09331 [hep-ph]
- AKC, Mitter, Pawłowski, Strodthoff,  $N_f = 2$  Vacuum QCD, in preparation
- AKC, Mitter, Pawłowski, Strodthoff,  $T > 0$  Yang-Mills, in preparation

November 9, 2016

# QCD phase diagram with functional methods

## fQCD-collaboration:

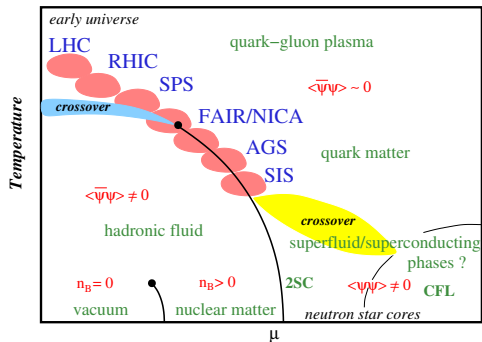
J. Braun, L. Corell, AKC, L. Fister, W. J. Fu, M. Leonhardt, M. Mitter,  
J. M. Pawłowski, M. Pospiech, F. Rennecke, N. Strodthoff, N. Wink, ...

## This talk:

- Vacuum Yang-Mills theory
- Preliminary  $T > 0$  results

## Aim:

- Qualitative understanding
- Quantitative precision



Schaefer and Wagner,  
 Prog.Part.Nucl.Phys. 62 (2009) 381

# QCD phase diagram with functional methods

## fQCD-collaboration:

J. Braun, L. Corell, AKC, L. Fister, W. J. Fu, M. Leonhardt, M. Mitter,  
J. M. Pawłowski, M. Pospiech, F. Rennecke, N. Strodthoff, N. Wink, ...

## This talk:

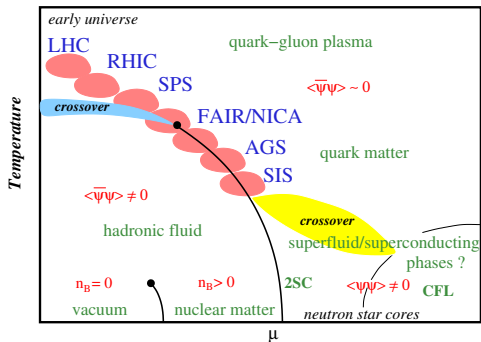
- Vacuum Yang-Mills theory
- Preliminary  $T > 0$  results

## Aim:

- Qualitative understanding
- Quantitative precision

## Motivation:

- No sign problem
- Understanding of confinement
- Prerequisite for QCD



Schaefer and Wagner,  
 Prog.Part.Nucl.Phys. 62 (2009) 381

# QCD phase diagram with functional methods

## fQCD-collaboration:

J. Braun, L. Corell, AKC, L. Fister, W. J. Fu, M. Leonhardt, M. Mitter,  
J. M. Pawłowski, M. Pospiech, F. Rennecke, N. Strodthoff, N. Wink, ...

## This talk:

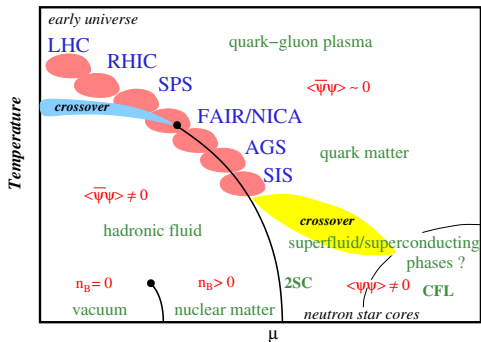
- Vacuum Yang-Mills theory
- Preliminary  $T > 0$  results

## Aim:

- Qualitative understanding
- Quantitative precision

## Motivation:

- No sign problem
- Understanding of confinement
- Prerequisite for QCD



Schaefer and Wagner,  
 Prog.Part.Nucl.Phys. 62 (2009) 381

# QCD phase diagram with functional methods

## fQCD-collaboration:

J. Braun, L. Corell, AKC, L. Fister, W. J. Fu, M. Leonhardt, M. Mitter,  
J. M. Pawłowski, M. Pospiech, F. Rennecke, N. Strodthoff, N. Wink, ...

## This talk:

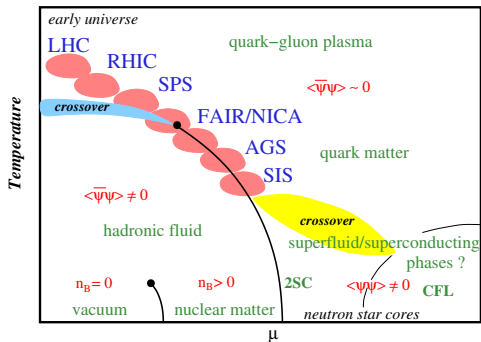
- Vacuum Yang-Mills theory
- Preliminary  $T > 0$  results

## Aim:

- Qualitative understanding
- Quantitative precision

## Motivation:

- No sign problem
- Understanding of confinement
- Prerequisite for QCD



Schaefer and Wagner,  
 Prog.Part.Nucl.Phys. 62 (2009) 381

# QCD from the functional renormalization group

- Only perturbative QCD input
  - $\alpha_S(\mu = \mathcal{O}(10) \text{ GeV})$
  - $m_q(\mu = \mathcal{O}(10) \text{ GeV})$
- Wetterich equation with initial condition  $S[\Phi] = \Gamma_\Lambda[\Phi]$
- Effective action  $\Gamma[\Phi] = \lim_{k \rightarrow 0} \Gamma_k[\Phi]$
- Exact equation
- $\partial_t$ : integration of momentum shells controlled by regulator
- Full field-dependent equation with  $(\Gamma^{(2)}[\Phi])^{-1}$  on rhs

# QCD from the functional renormalization group

- Only perturbative QCD input

- $\alpha_S(\mu = \mathcal{O}(10) \text{ GeV})$
- $m_q(\mu = \mathcal{O}(10) \text{ GeV})$

- Wetterich equation with initial condition  $S[\Phi] = \Gamma_\Lambda[\Phi]$

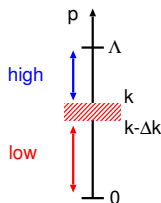
$$\partial_t \Gamma_k[\phi] = +\frac{1}{2} \left( \text{Diagram 1} - \text{Diagram 2} - \text{Diagram 3} \right)$$

- Effective action  $\Gamma[\Phi] = \lim_{k \rightarrow 0} \Gamma_k[\Phi]$

- Exact equation

- $\partial_t$ : integration of momentum shells controlled by regulator

- Full field-dependent equation with  $(\Gamma^{(2)}[\Phi])^{-1}$  on rhs



# QCD from the functional renormalization group

- Only perturbative QCD input

- $\alpha_S(\mu = \mathcal{O}(10) \text{ GeV})$
- $m_q(\mu = \mathcal{O}(10) \text{ GeV})$

- Wetterich equation with initial condition  $S[\Phi] = \Gamma_\Lambda[\Phi]$

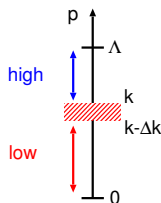
$$\partial_t \Gamma_k[\phi] = +\frac{1}{2} \left( \text{Diagram 1} - \text{Diagram 2} - \text{Diagram 3} \right)$$

- Effective action  $\Gamma[\Phi] = \lim_{k \rightarrow 0} \Gamma_k[\Phi]$

- Exact equation

- $\partial_t$ : integration of momentum shells controlled by regulator

- Full field-dependent equation with  $(\Gamma^{(2)}[\Phi])^{-1}$  on rhs





## Vertex expansion

- Approximation necessary – vertex expansion:

$$\Gamma[\Phi] = \sum_n \int_{p_1, \dots, p_{n-1}} \Gamma_{\Phi_1 \dots \Phi_n}^{(n)}(p_1, \dots, p_{n-1}) \Phi^1(p_1) \cdots \Phi^n(-p_1 - \dots - p_{n-1})$$

- Wanted: “apparent convergence” of  $\Gamma[\Phi] = \lim_{k \rightarrow 0} \Gamma_k[\Phi]$
- Current state-of-the-art truncation:

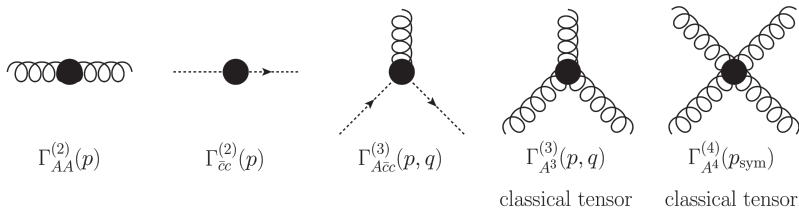
- Functional derivatives of  $\Gamma_k[\Phi]$  with respect to fields yield equations

# Vertex expansion

- Approximation necessary – vertex expansion:

$$\Gamma[\Phi] = \sum_n \int_{p_1, \dots, p_{n-1}} \Gamma_{\Phi_1 \dots \Phi_n}^{(n)}(p_1, \dots, p_{n-1}) \Phi^1(p_1) \cdots \Phi^n(-p_1 - \dots - p_{n-1})$$

- Wanted: “apparent convergence” of  $\Gamma[\Phi] = \lim_{k \rightarrow 0} \Gamma_k[\Phi]$
- Current state-of-the-start truncation:



- Functional derivatives of  $\Gamma_k[\Phi]$  with respect to fields yield equations

# Truncation – closed set of equations

$$\partial_t \text{---}^{-1} = \text{---} \text{---} \text{---} + \text{---} \text{---} \text{---}$$

$$\partial_t \text{---}^{-1} = \text{---} \text{---} \text{---} - 2 \text{---} \text{---} \text{---} - \frac{1}{2} \text{---} \text{---} \text{---}$$

$$\partial_t \text{---} = - \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} + \text{perm.}$$

$$\partial_t \text{---} = - \text{---} \text{---} \text{---} + 2 \text{---} \text{---} \text{---} + \text{---} \text{---} \text{---} + \text{perm.}$$

$$\partial_t \text{---} = + \text{---} \text{---} \text{---} + \text{---} \text{---} \text{---} - 2 \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} + \text{perm.}$$

# Truncation – closed set of equations

$$\partial_t \text{---}^{-1} = \text{---} \otimes \text{---} + \text{---} \otimes \text{---}$$

momentum dependent  
coupled **tensor equations**

$$\partial_t \text{---}^{-1} = \text{---} \otimes \text{---} - 2 \text{---} \otimes \text{---} - \frac{1}{2} \text{---} \otimes \text{---}$$

$$\partial_t \text{---} = - \text{---} \otimes \text{---} - \text{---} \otimes \text{---} + \text{perm.}$$

$$\partial_t \text{---} = - \text{---} \otimes \text{---} + 2 \text{---} \otimes \text{---} + \text{---} \otimes \text{---} + \text{perm.}$$

$$\partial_t \text{---} = + \text{---} \otimes \text{---} + \text{---} \otimes \text{---} - 2 \text{---} \otimes \text{---} - \text{---} \otimes \text{---} + \text{perm.}$$

# Truncation – closed set of equations

$$\partial_t \text{---}^{-1} = \text{---} \text{---} \text{---} + \text{---} \text{---} \text{---}$$

momentum dependent  
coupled **tensor equations**

$$\partial_t \text{---}^{-1} = \text{---} \text{---} \text{---} - 2 \text{---} \text{---} \text{---} - \frac{1}{2} \text{---} \text{---} \text{---}$$

tracing necessary

$$\partial_t \text{---} \text{---} \text{---} = - \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} + \text{perm.}$$

$$\partial_t \text{---} \text{---} \text{---} = - \text{---} \text{---} \text{---} + 2 \text{---} \text{---} \text{---} + \text{---} \text{---} \text{---} + \text{perm.}$$

$$\partial_t \text{---} \text{---} \text{---} = + \text{---} \text{---} \text{---} + \text{---} \text{---} \text{---} - 2 \text{---} \text{---} \text{---} - \text{---} \text{---} \text{---} + \text{perm.}$$

# FormTracer – Mathematica tracing package using FORM

- Mathematica: very powerful, flexible and **convenient**
- FORM: very **fast** and **efficient**

**FormTracer** uses FORM while it keeps the usability of Mathematica:

- Lorentz/Dirac traces in arbitrary dimensions
- Arbitrary number of group product spaces
- Intuitive, easy-to-use and highly customizable Mathematica frontend
- Support for finite temperature/density applications
- Support for FORM's optimization algorithm
- Convenient installation and update procedure within Mathematica:

**Preprint: AKC, Mitter, Strodthoff; arXiv:1610.09331 [hep-ph]**

**Open source:** <https://github.com/FormTracer/FormTracer>

# FormTracer – installation and usage

FormTracer.nb - Wolfram Mathematica 11.0

File Edit Insert Format Cell Graphics Evaluation Palettes Window Help

## Installing

```
Import["https://raw.githubusercontent.com/FormTracer/FormTracer/master/src/FormTracerInstaller.m"]
```

## Tracing

### Space-Time

#### Define syntax for space-time

```
DefineLorentzTensors[ $\delta[\mu, \nu]$  (*Kronecker delta*),  $\text{vec}[p, \mu]$  (*vector*),  $p \cdot q$  (*inner product*)];
```

#### Take traces:

```
FormTrace[ $\text{vec}[p + 2 r, \mu] \delta[\mu, \nu] \text{vec}[s, \nu]$ ]
FormTrace[ $\delta[\alpha, \nu] (\delta[\nu, \rho] + \delta[\nu, \rho] \delta[\sigma, \sigma]) \delta[\rho, \alpha]$ ]
FormTrace[ $\delta[1, \nu] \text{vec}[s, \nu]$ ]
s . (p + 2 r)
20
vec[s, 1]
```

**AKC, Mitter, Strodthoff; arXiv:1610.09331 [hep-ph]**

# Truncation – closed set of equations

$$\partial_t \text{---}^{-1} = \text{---} \otimes \text{---} + \text{---} \otimes \text{---}$$

momentum dependent  
coupled tensor equations

$$\partial_t \text{---}^{-1} = \text{---} \otimes \text{---} - 2 \text{---} \otimes \text{---} - \frac{1}{2} \text{---} \otimes \text{---}$$

tracing necessary

$$\partial_t \text{---} = - \text{---} \otimes \text{---} - \text{---} \otimes \text{---} + \text{perm.}$$

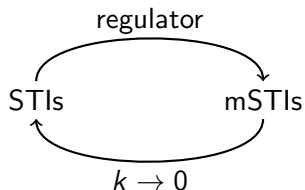
$$\partial_t \text{---} = - \text{---} \otimes \text{---} + 2 \text{---} \otimes \text{---} + \text{---} \otimes \text{---} + \text{perm.}$$

$$\partial_t \text{---} = + \text{---} \otimes \text{---} + \text{---} \otimes \text{---} - 2 \text{---} \otimes \text{---} - \text{---} \otimes \text{---} + \text{perm.}$$



# Regulator breaks BRST symmetry

- Breaking BRST symmetry  $\rightarrow$  modified STIs
- mSTIs reduce to STIs at  $k = 0$
- $\implies$  solve mSTIs to get initial action at  $k = \Lambda$
- More practical solution: choose  $\Gamma_\Lambda \approx S$  such that STIs are fulfilled  $k = 0$



$$\alpha_{A\bar{c}c}(p) = \frac{\alpha(\mu)}{4\pi} \frac{Z_{A\bar{c}c}^2(p)}{Z_A(p) Z_c^2(p)}$$

Select

$$Z_{A\bar{c}c}^{k=\Lambda}(p) = \text{const.}$$

$$\alpha_{A^3}(p) = \frac{\alpha(\mu)}{4\pi} \frac{Z_{A^3}^2(p)}{Z_A^3(p)}$$

$$Z_{A^3}^{k=\Lambda}(p) = \text{const.}$$

$$\alpha_{A^4}(p) = \frac{\alpha(\mu)}{4\pi} \frac{Z_{A^4}(p)}{Z_A^2(p)}$$

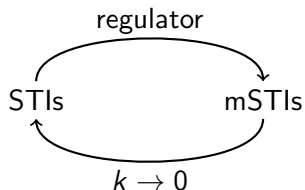
$$Z_{A^4}^{k=\Lambda}(p) = \text{const.}$$

such that

$$\alpha_{A\bar{c}c}(\mu) = \alpha_{A^3}(\mu) = \alpha_{A^4}(\mu)$$

# Regulator breaks BRST symmetry

- Breaking BRST symmetry  $\rightarrow$  modified STIs
- mSTIs reduce to STIs at  $k = 0$
- $\implies$  solve mSTIs to get initial action at  $k = \Lambda$
- More practical solution: choose  $\Gamma_\Lambda \approx S$  such that STIs are fulfilled  $k = 0$



$$\alpha_{A\bar{c}c}(p) = \frac{\alpha(\mu)}{4\pi} \frac{Z_{A\bar{c}c}^2(p)}{Z_A(p) Z_c^2(p)}$$

$$\alpha_{A^3}(p) = \frac{\alpha(\mu)}{4\pi} \frac{Z_{A^3}^2(p)}{Z_A^3(p)}$$

$$\alpha_{A^4}(p) = \frac{\alpha(\mu)}{4\pi} \frac{Z_{A^4}(p)}{Z_A^2(p)}$$

Select

$$Z_{A\bar{c}c}^{k=\Lambda}(p) = \text{const.}$$

$$Z_{A^3}^{k=\Lambda}(p) = \text{const.}$$

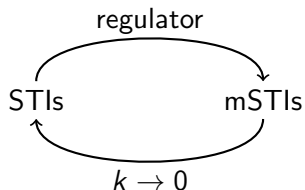
$$Z_{A^4}^{k=\Lambda}(p) = \text{const.}$$

such that

$$\alpha_{A\bar{c}c}(\mu) = \alpha_{A^3}(\mu) = \alpha_{A^4}(\mu)$$

# Regulator breaks BRST symmetry

- Breaking BRST symmetry  $\rightarrow$  modified STIs
- mSTIs reduce to STIs at  $k = 0$
- $\implies$  solve mSTIs to get initial action at  $k = \Lambda$
- More practical solution: choose  $\Gamma_\Lambda \approx S$  such that STIs are fulfilled  $k = 0$



$$\alpha_{A\bar{c}c}(p) = \frac{\alpha(\mu)}{4\pi} \frac{Z_{A\bar{c}c}^2(p)}{Z_A(p) Z_c^2(p)}$$

$$\alpha_{A^3}(p) = \frac{\alpha(\mu)}{4\pi} \frac{Z_{A^3}^2(p)}{Z_A^3(p)}$$

$$\alpha_{A^4}(p) = \frac{\alpha(\mu)}{4\pi} \frac{Z_{A^4}(p)}{Z_A^2(p)}$$

Select

$$Z_{A\bar{c}c}^{k=\Lambda}(p) = \text{const.}$$

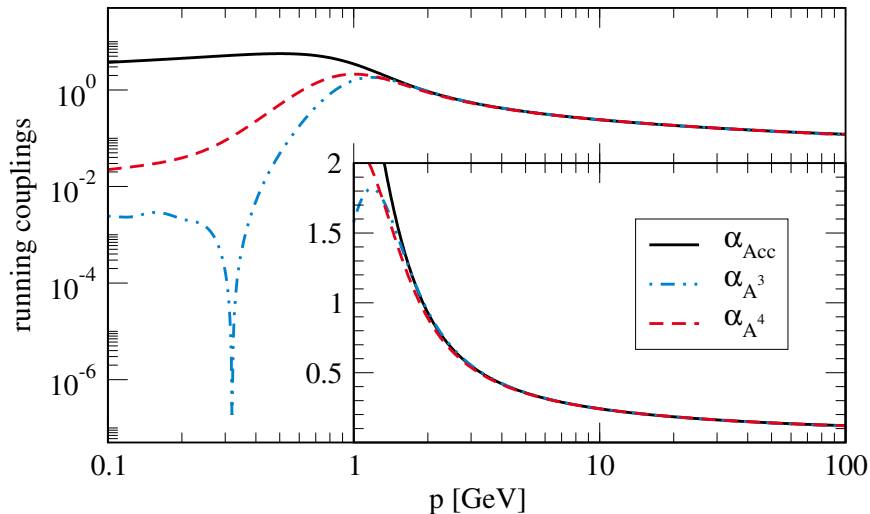
$$Z_{A^3}^{k=\Lambda}(p) = \text{const.}$$

$$Z_{A^4}^{k=\Lambda}(p) = \text{const.}$$

such that

$$\alpha_{A\bar{c}c}(\mu) = \alpha_{A^3}(\mu) = \alpha_{A^4}(\mu)$$

# Running couplings (scaling solution)



AKC, Fister, Mitter, Pawłowski, Strodthoff, 2016

# Gluon mass gap

Scaling solution

$$\lim_{p \rightarrow 0} Z_c(p^2) \propto (p^2)^\kappa$$

$$\lim_{p \rightarrow 0} Z_A(p^2) \propto (p^2)^{-2\kappa}$$

Decoupling solution

$$\lim_{p \rightarrow 0} Z_c(p^2) \propto 1$$

$$\lim_{p \rightarrow 0} Z_A(p^2) \propto (p^2)^{-1}$$

- Landau Gauge gluon STI requires longitudinally mass term to vanish:

$$p_\mu \left( [\Gamma_{AA,L}^{(2)}]_{\mu\nu}^{ab}(p) - [S_{AA,L}^{(2)}]_{\mu\nu}^{ab}(p) \right) = 0$$

- Splitting between longitudinal and transverse mass term necessary
- Splitting occurs "naturally" for scaling solution
- Decoupling solution requires irregular vertices, e.g. a pole in the longitudinal sector
- Unphysical mass parameter at cutoff:  $[\Gamma_{AA}^{(2)}]_{\mu\nu}^{ab}(k = \Lambda, p) = p^2 + m_\Lambda^2$ ,
- $m_\Lambda^2 \propto c \cdot \Lambda^2$  with  $c \ll 1$ ; parameter can be uniquely determined

# Gluon mass gap

Scaling solution

$$\lim_{p \rightarrow 0} Z_c(p^2) \propto (p^2)^\kappa$$

$$\lim_{p \rightarrow 0} Z_A(p^2) \propto (p^2)^{-2\kappa}$$

Decoupling solution

$$\lim_{p \rightarrow 0} Z_c(p^2) \propto 1$$

$$\lim_{p \rightarrow 0} Z_A(p^2) \propto (p^2)^{-1}$$

- **Landau Gauge** gluon STI requires longitudinally mass term to vanish:

$$p_\mu \left( [\Gamma_{AA,L}^{(2)}]_{\mu\nu}^{ab}(p) - [S_{AA,L}^{(2)}]_{\mu\nu}^{ab}(p) \right) = 0$$

- Splitting between longitudinal and transverse mass term necessary
- Splitting occurs "naturally" for scaling solution
- Decoupling solution requires irregular vertices, e.g. a pole in the longitudinal sector
- Unphysical mass parameter at cutoff:  $[\Gamma_{AA}^{(2)}]_{\mu\nu}^{ab}(k = \Lambda, p) = p^2 + m_\Lambda^2$ ,
- $m_\Lambda^2 \propto c \cdot \Lambda^2$  with  $c \ll 1$ ; parameter can be uniquely determined

# Gluon mass gap

Scaling solution

$$\lim_{p \rightarrow 0} Z_c(p^2) \propto (p^2)^\kappa$$

$$\lim_{p \rightarrow 0} Z_A(p^2) \propto (p^2)^{-2\kappa}$$

Decoupling solution

$$\lim_{p \rightarrow 0} Z_c(p^2) \propto 1$$

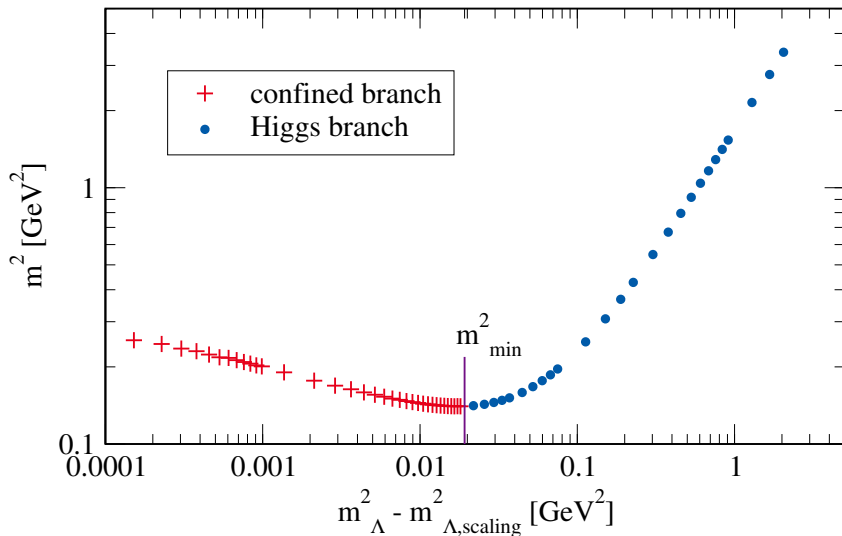
$$\lim_{p \rightarrow 0} Z_A(p^2) \propto (p^2)^{-1}$$

- **Landau Gauge** gluon STI requires longitudinally mass term to vanish:

$$p_\mu \left( [\Gamma_{AA,L}^{(2)}]_{\mu\nu}^{ab}(p) - [S_{AA,L}^{(2)}]_{\mu\nu}^{ab}(p) \right) = 0$$

- Splitting between longitudinal and transverse mass term necessary
- Splitting occurs "naturally" for scaling solution
- Decoupling solution requires irregular vertices, e.g. a pole in the longitudinal sector
- Unphysical mass parameter at cutoff:  $[\Gamma_{AA}^{(2)}]_{\mu\nu}^{ab}(k = \Lambda, p) = p^2 + m_\Lambda^2$ ,
- $m_\Lambda^2 \propto c \cdot \Lambda^2$  with  $c \ll 1$ ; **parameter can be uniquely determined**

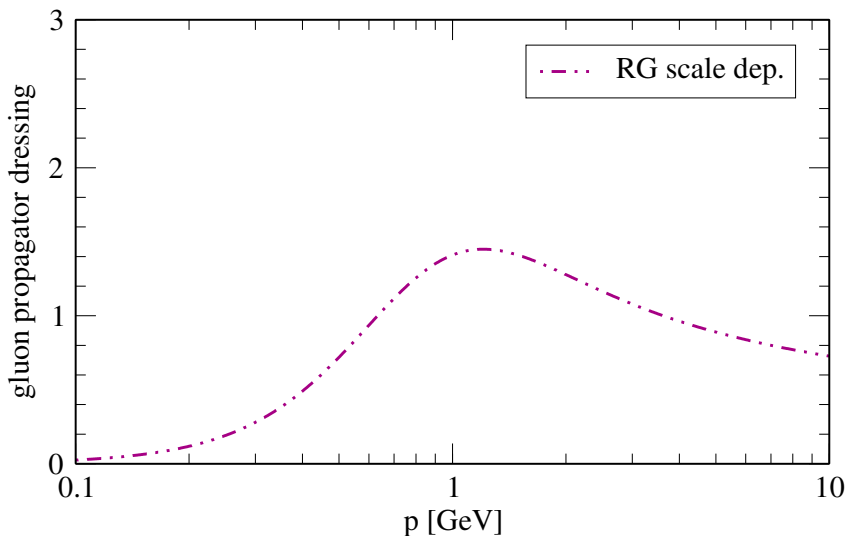
# Dynamical mass generation



AKC, Fister, Mitter, Pawłowski, Strodthoff, 2016

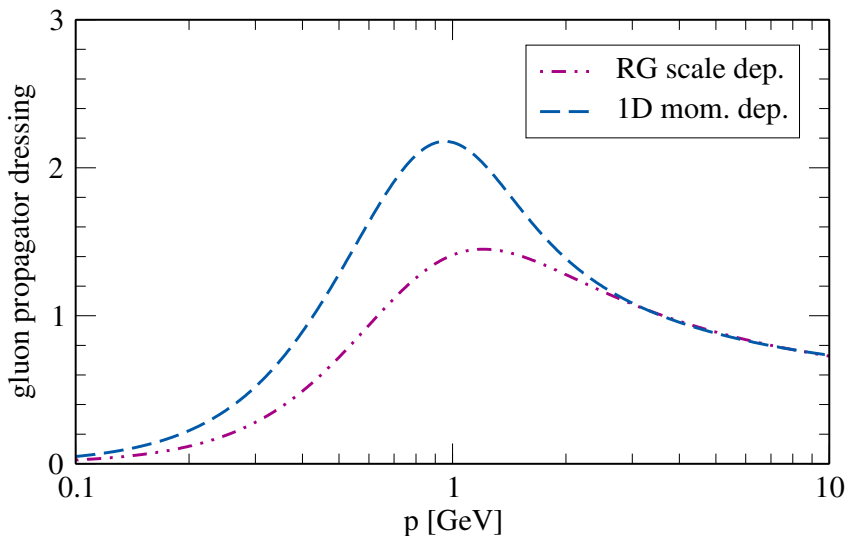


# Truncation dependence of the gluon propagator



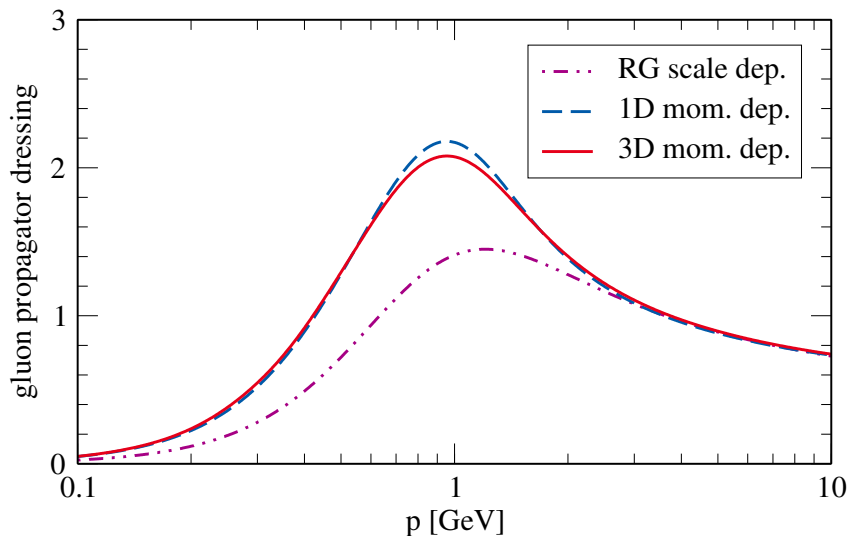
AKC, Fister, Mitter, Pawłowski, Strodthoff, 2016

# Truncation dependence of the gluon propagator



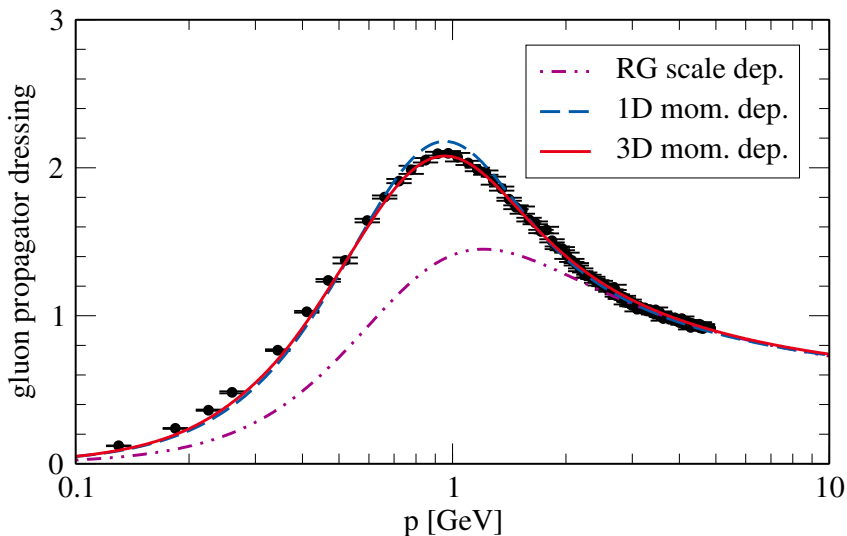
AKC, Fister, Mitter, Pawłowski, Strodthoff, 2016

# Truncation dependence of the gluon propagator



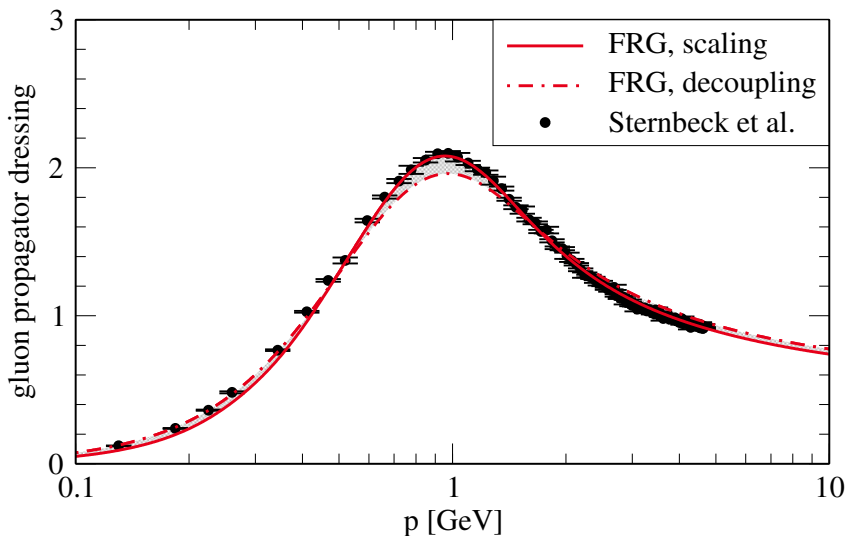
AKC, Fister, Mitter, Pawłowski, Strodthoff, 2016

# Truncation dependence of the gluon propagator



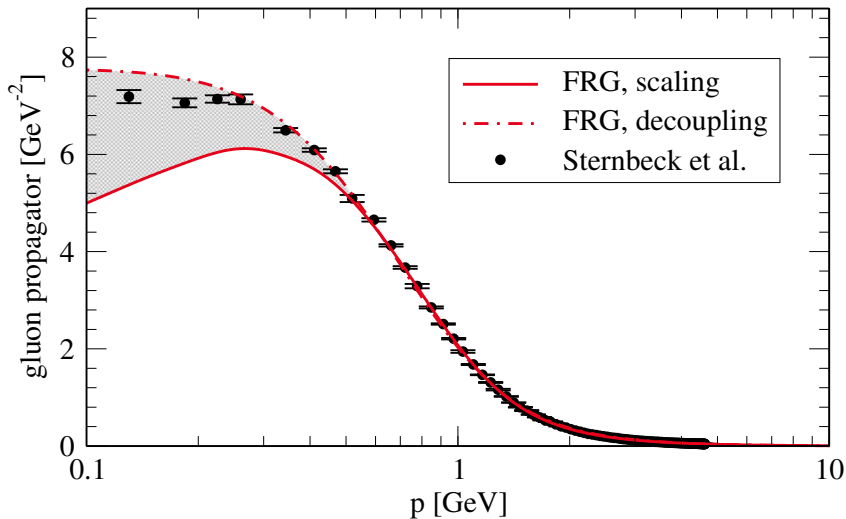
AKC, Fister, Mitter, Pawłowski, Strodthoff, 2016; **Lattice:** Sternbeck et al. 2006

# Gluon propagator dressing



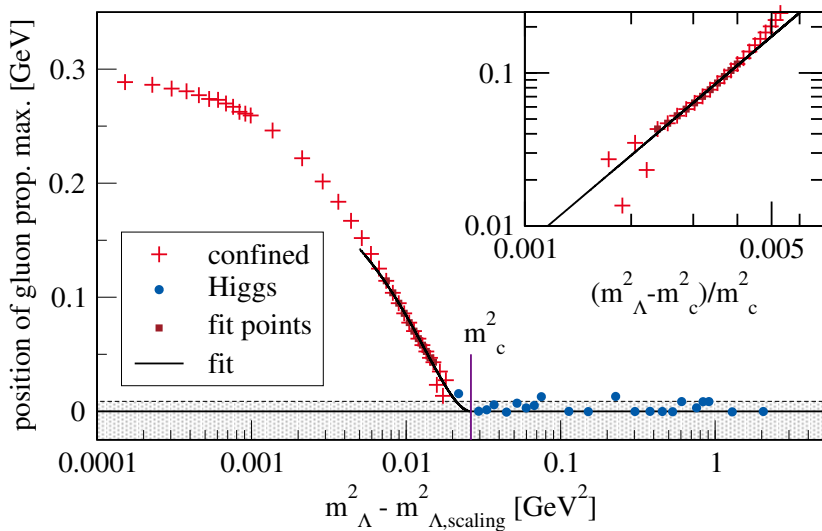
AKC, Fister, Mitter, Pawłowski, Strodthoff, 2016; **Lattice**: Sternbeck et al. 2006

# Gluon propagator



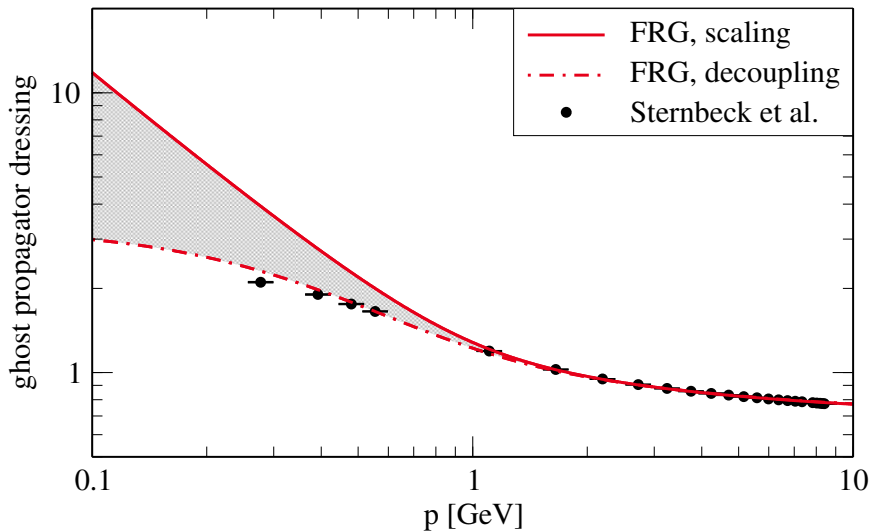
AKC, Fister, Mitter, Pawłowski, Strodthoff, 2016; **Lattice**: Sternbeck et al. 2006

# Gluon propagator maximum over UV mass parameter



AKC, Fister, Mitter, Pawłowski, Strodthoff, 2016

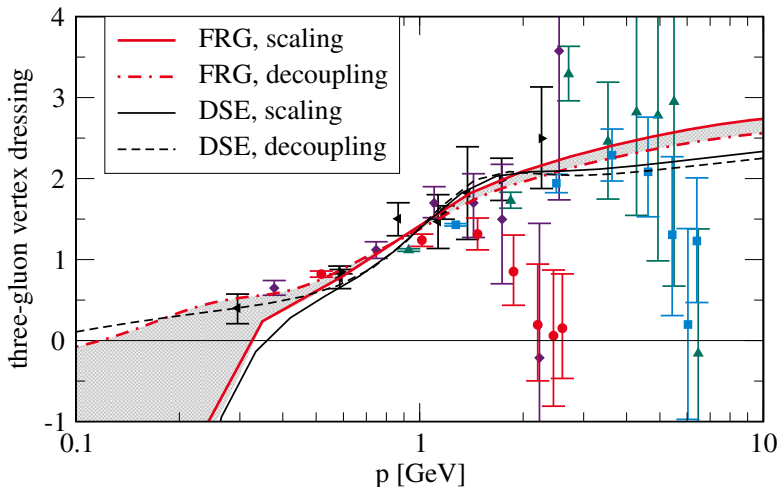
# Ghost propagator dressing



AKC, Fister, Mitter, Pawłowski, Strodthoff, 2016



# Three-gluon vertex dressing (symmetric point)



- Zero crossing between 0.1 GeV to 0.33 GeV

AKC, Fister, Mitter, Pawłowski, Strodthoff, 2016

# Finite temperature

## Going to finite temperature:

- Introduce Matsubara frequencies:

$$\int \frac{d^4 p}{(2\pi)^4} \rightarrow T \sum_{\omega_n} \int \frac{d^3 p}{(2\pi)^3}$$

- Thermal Debye mass
- Same parameter-free truncation as in vacuum YM
- Upcoming: full splitting of magnetic and electric components

Splitting of propagators only: Fister, Pawłowski, 2011

$$P_{\mu\nu}^T(p) = (1 - \delta_{0\mu})(1 - \delta_{0\nu}) \left( \delta_{\mu\nu} - \frac{p_\mu p_\nu}{p^2} \right) \quad P_{\mu\nu}^L(p) = \left( \delta_{\mu\nu} - \frac{p_\mu p_\nu}{p^2} \right) - P_{\mu\nu}^T(p)$$

- Also upcoming: nonzero Matsubara modes

## Following results are preliminary and based on

- AKC, Mitter, Pawłowski, Strodthoff,  $T > 0$  Yang-Mills, in preparation

# Finite temperature

## Going to finite temperature:

- Introduce Matsubara frequencies:

$$\int \frac{d^4 p}{(2\pi)^4} \rightarrow T \sum_{\omega_n} \int \frac{d^3 p}{(2\pi)^3}$$

- Thermal Debye mass
- Same parameter-free truncation as in vacuum YM
- Upcoming: full splitting of magnetic and electric components

Splitting of propagators only: Fister, Pawłowski, 2011

$$P_{\mu\nu}^T(p) = (1 - \delta_{0\mu})(1 - \delta_{0\nu}) \left( \delta_{\mu\nu} - \frac{p_\mu p_\nu}{\vec{p}^2} \right) \quad P_{\mu\nu}^L(p) = \left( \delta_{\mu\nu} - \frac{p_\mu p_\nu}{p^2} \right) - P_{\mu\nu}^T(p)$$

- Also upcoming: nonzero Matsubara modes

Following results are preliminary and based on

- AKC, Mitter, Pawłowski, Strodthoff,  $T > 0$  Yang-Mills, in preparation

# Finite temperature

## Going to finite temperature:

- Introduce Matsubara frequencies:

$$\int \frac{d^4 p}{(2\pi)^4} \rightarrow T \sum_{\omega_n} \int \frac{d^3 p}{(2\pi)^3}$$

- Thermal Debye mass
- Same parameter-free truncation as in vacuum YM
- Upcoming: full splitting of magnetic and electric components

Splitting of propagators only: Fister, Pawłowski, 2011

$$P_{\mu\nu}^T(p) = (1 - \delta_{0\mu})(1 - \delta_{0\nu}) \left( \delta_{\mu\nu} - \frac{p_\mu p_\nu}{\vec{p}^2} \right) \quad P_{\mu\nu}^L(p) = \left( \delta_{\mu\nu} - \frac{p_\mu p_\nu}{p^2} \right) - P_{\mu\nu}^T(p)$$

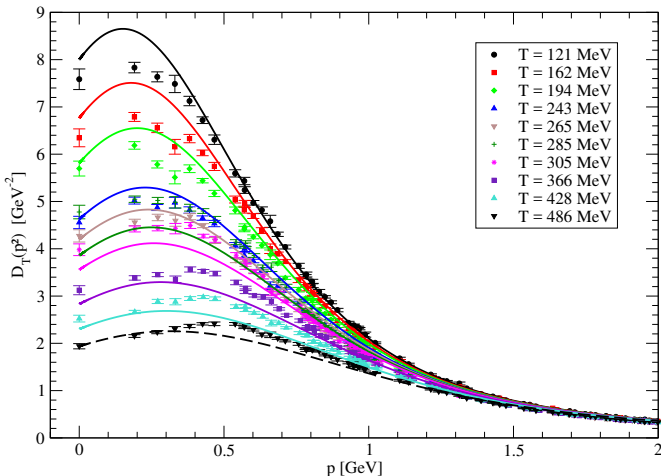
- Also upcoming: nonzero Matsubara modes

## Following results are preliminary and based on

- AKC, Mitter, Pawłowski, Strodthoff,  $T > 0$  Yang-Mills, in preparation

# Temperature dependence of the gluon propagator

Magnetic component compared to averaged components from FRG:



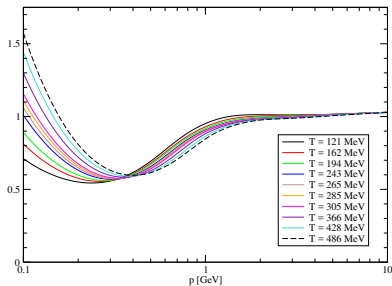
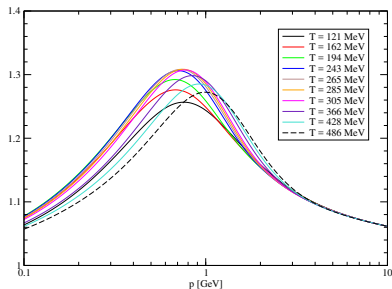
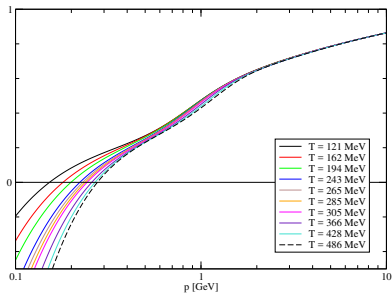
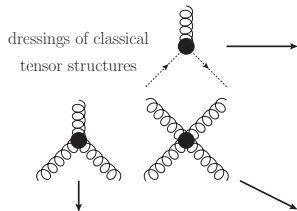
Zeroth mode results:

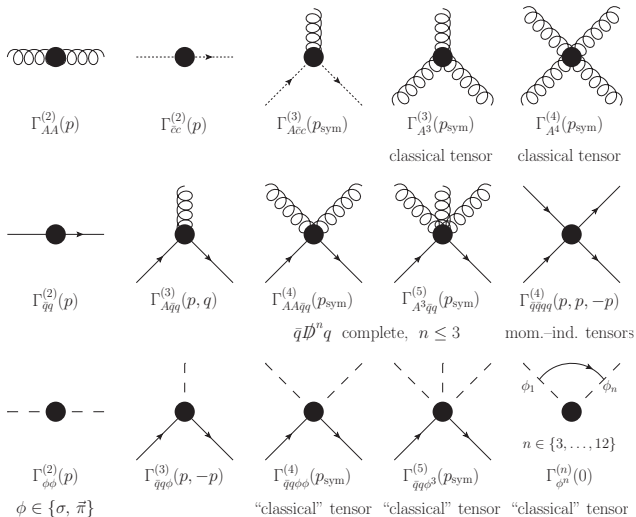
Lattice: Silva, Oliveira,  
Bicudo, Cardoso, 2013

FRG: AKC, Mitter,  
Pawlowski, Strodthoff,  
preliminary

# Temperature dependence of vertices

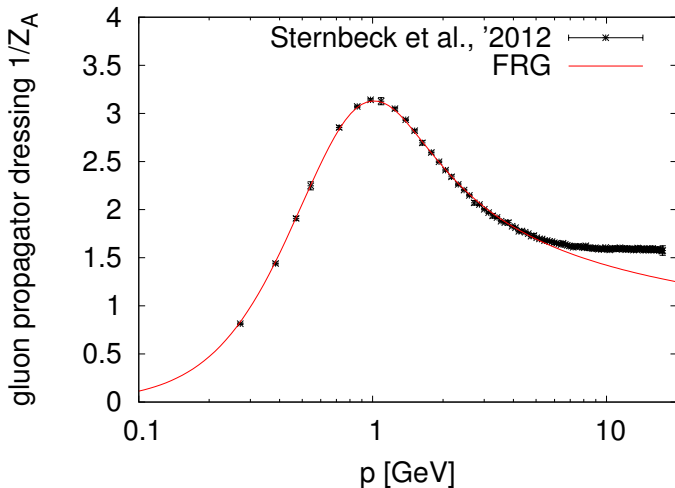
AKC, Mitter, Pawłowski, Strodthoff, preliminary



A glimpse at unquenched  $N_f = 2$  QCD

AKC, Mitter, Pawłowski, Strodthoff, in preparation

## Preliminary unquenched gluon propagator

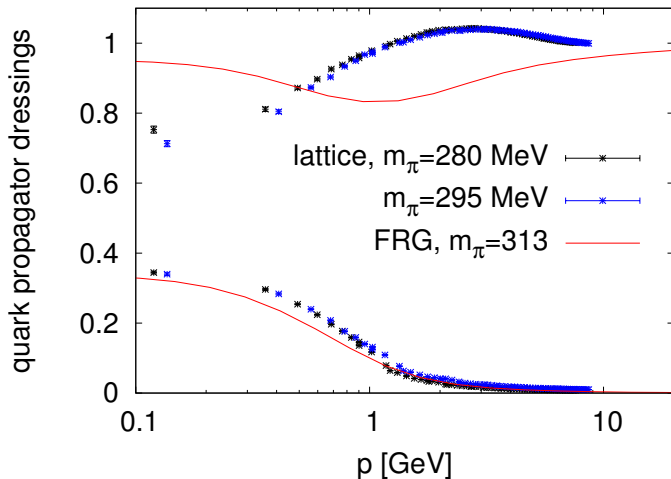


FRG: AKC, Mitter, Pawłowski, Strodthoff, in preparation

Lattice: A. Sternbeck, K. Maltman, M. Müller-Preussker, L. von Smekal, PoS LATTICE2012 (2012) 243



## Preliminary unquenched quark propagator



FRG: AKC, Mitter, Pawłowski, Strodthoff, in preparation

Lattice: Oliveira, Kızılersu, Silva, Skullerud, Sternbeck, Williams, arXiv:1605.09632 [hep-lat]

## Conclusion

- FRG first principal approach to QCD, complementary to lattice QCD
- Big numerical effort  $\rightarrow$  tools like FormTracer necessary
- BRST symmetry is broken by regulator, proper care needs to be taken
- STI consistent solution computed
- Evidence for dynamical mass generation
- Very good agreement with lattice results

## Outlook

- Unquenched  $N_f = 2$  QCD, in preparation
- $T > 0$  YM with splitting of el. and mag. components, in preparation
- Bound states (Bethe-Salpeter eq.), decay widths, ...
- Nonzero Matsubara modes, gluon spectral function, ...

**Thank you for your attention!**

## Conclusion

- FRG first principal approach to QCD, complementary to lattice QCD
- Big numerical effort  $\rightarrow$  tools like FormTracer necessary
- BRST symmetry is broken by regulator, proper care needs to be taken
- STI consistent solution computed
- Evidence for dynamical mass generation
- Very good agreement with lattice results

## Outlook

- Unquenched  $N_f = 2$  QCD, in preparation
- $T > 0$  YM with splitting of el. and mag. components, in preparation
- Bound states (Bethe-Salpeter eq.), decay widths, ...
- Nonzero Matsubara modes, gluon spectral function, ...

Thank you for your attention!

## Conclusion

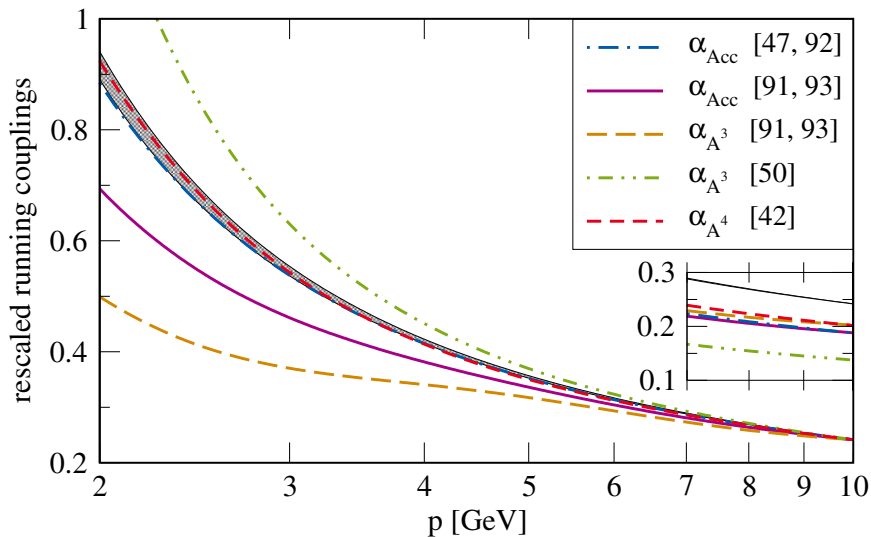
- FRG first principal approach to QCD, complementary to lattice QCD
- Big numerical effort  $\rightarrow$  tools like FormTracer necessary
- BRST symmetry is broken by regulator, proper care needs to be taken
- STI consistent solution computed
- Evidence for dynamical mass generation
- Very good agreement with lattice results

## Outlook

- Unquenched  $N_f = 2$  QCD, in preparation
- $T > 0$  YM with splitting of el. and mag. components, in preparation
- Bound states (Bethe-Salpeter eq.), decay widths, ...
- Nonzero Matsubara modes, gluon spectral function, ...

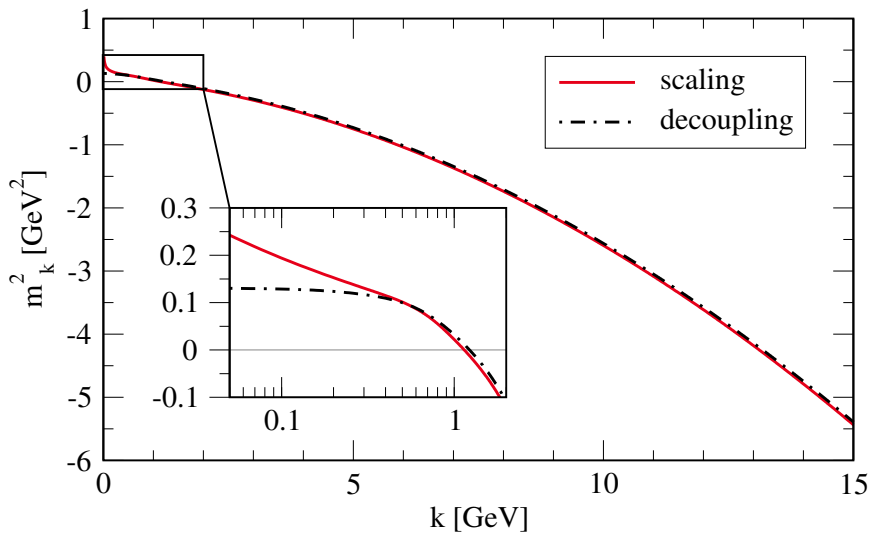
**Thank you for your attention!**

## Running couplings in comparison with DSE results



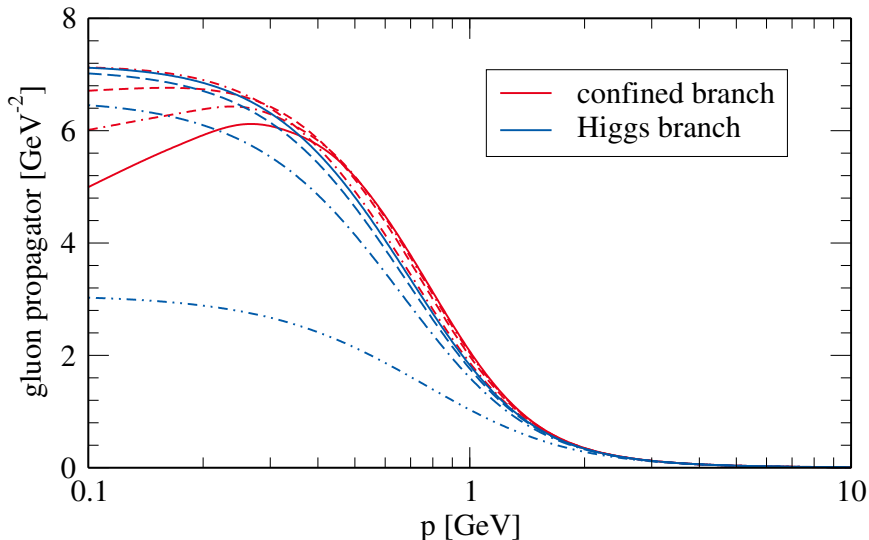
AKC, Fister, Mitter, Pawłowski, Strodthoff, 2016

## Running of the gluon mass parameter



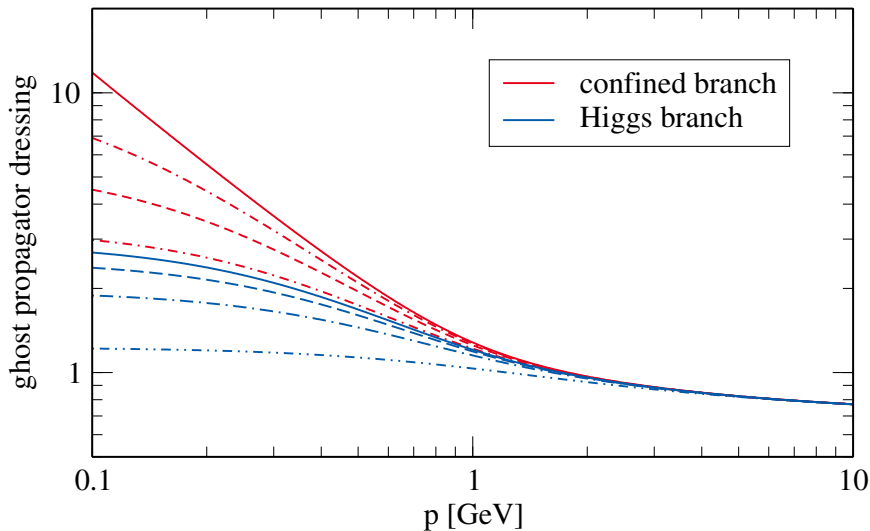
AKC, Fister, Mitter, Pawłowski, Strodthoff, 2016

## Gluon propagator



AKC, Fister, Mitter, Pawłowski, Strodthoff, 2016

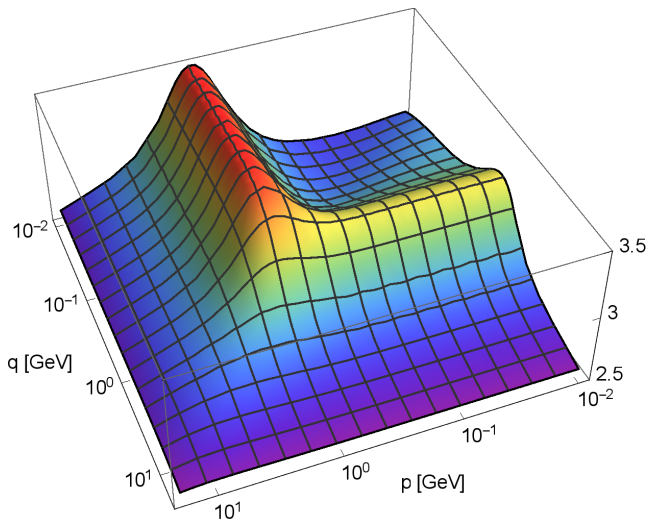
# Ghost propagator dressing



AKC, Fister, Mitter, Pawłowski, Strodthoff, 2016

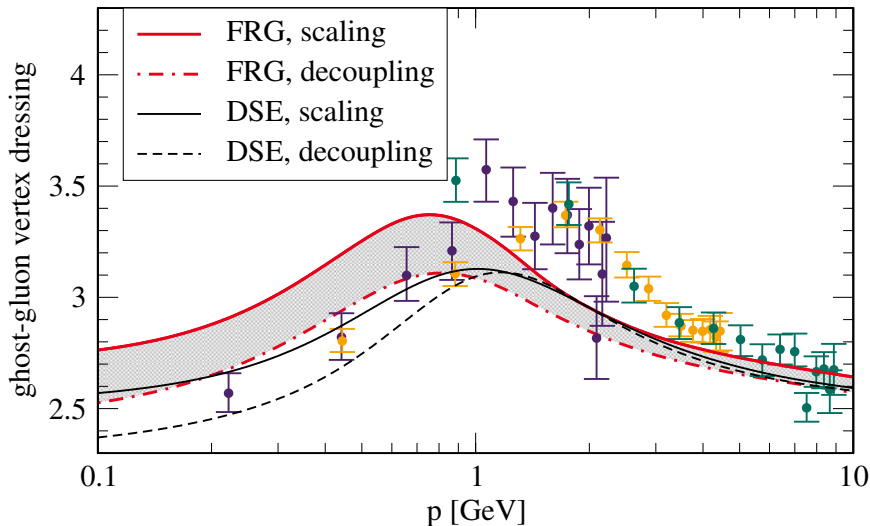


# Momentum dependence of the ghost-gluon vertex



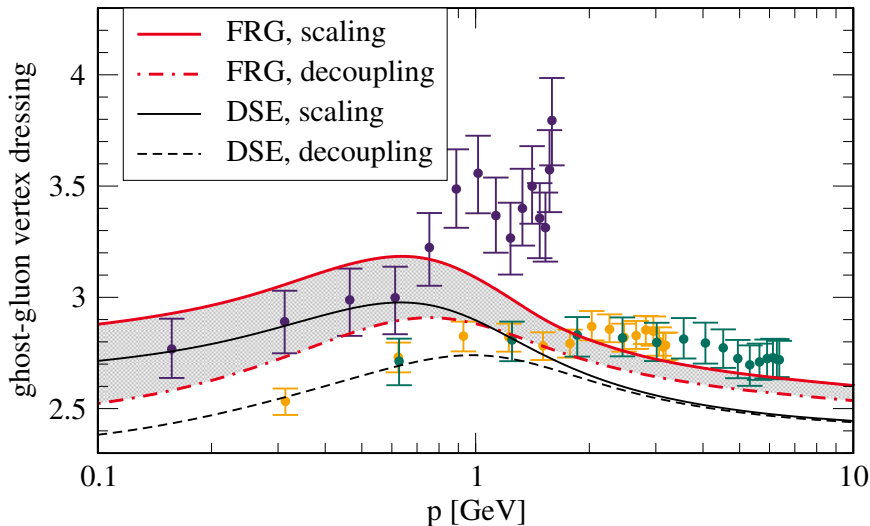
AKC, Fister, Mitter, Pawłowski, Strodthoff, 2016

# Ghost-gluon vertex at the symmetric point



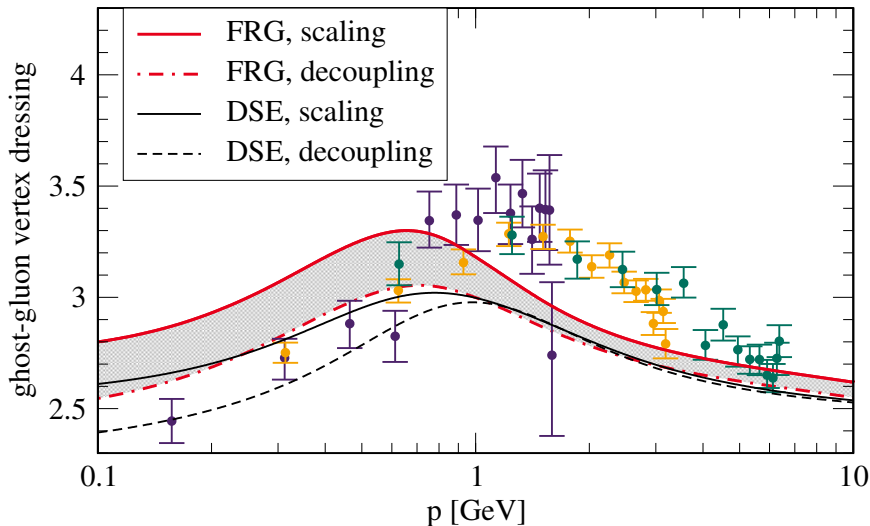
AKC, Fister, Mitter, Pawłowski, Strodthoff, 2016

# Ghost-gluon vertex with vanishing gluon momentum



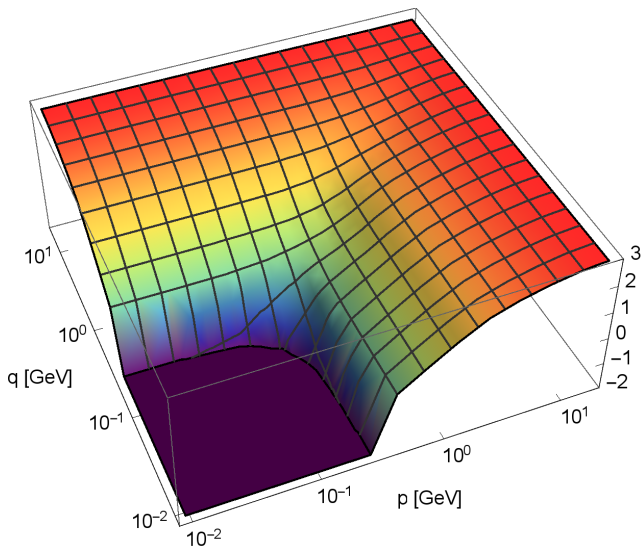
AKC, Fister, Mitter, Pawlowski, Strodthoff, 2016

# Ghost-gluon vertex with orthogonal momenta



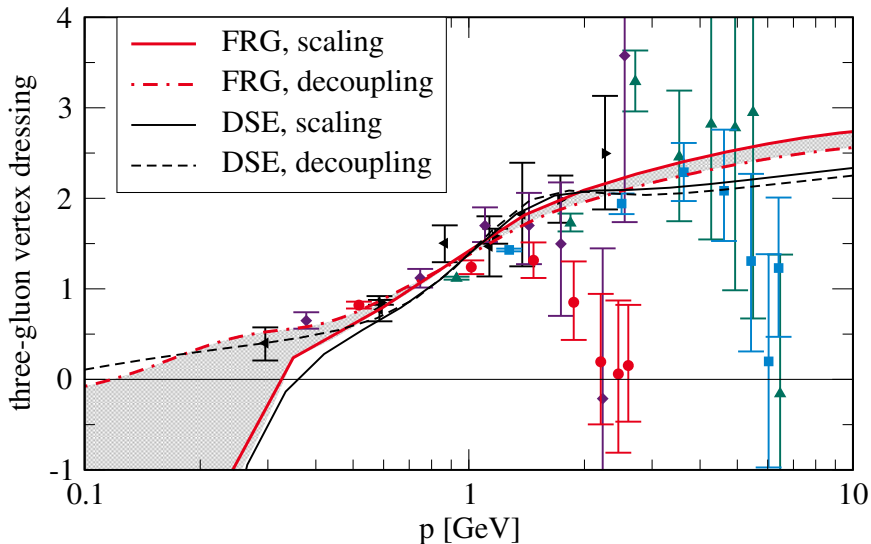
AKC, Fister, Mitter, Pawłowski, Strodthoff, 2016

# Momentum dependence of the three-gluon vertex



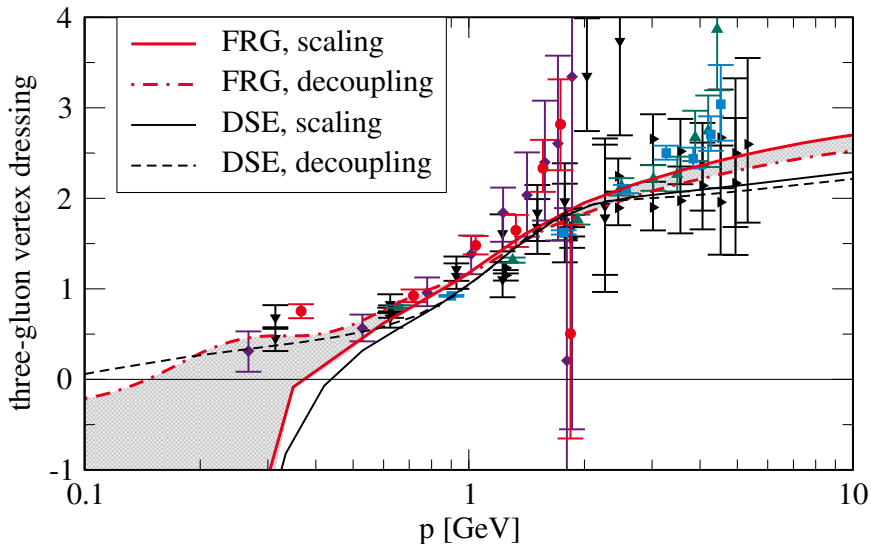
AKC, Fister, Mitter, Pawłowski, Strodthoff, 2016

# Three-gluon vertex at the symmetric point



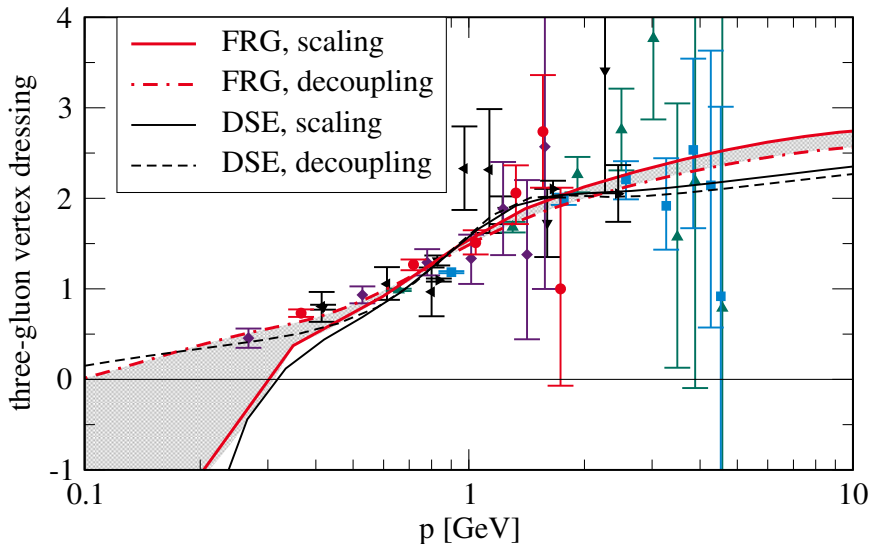
AKC, Fister, Mitter, Pawłowski, Strodthoff, 2016

# Three-gluon vertex with vanishing gluon momentum



AKC, Fister, Mitter, Pawłowski, Strodthoff, 2016

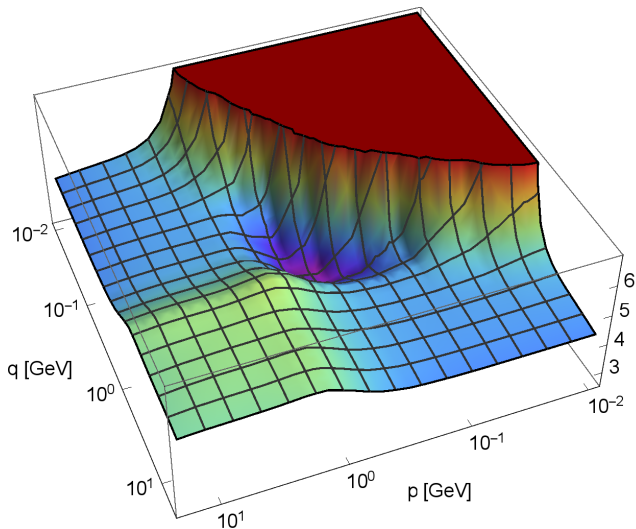
# Three-gluon vertex with orthogonal momenta



AKC, Fister, Mitter, Pawlowski, Strodthoff, 2016

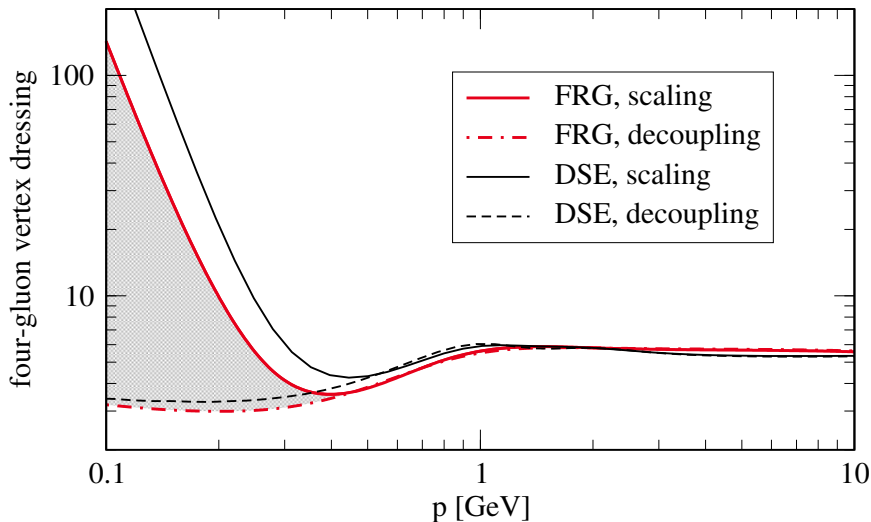


# Momentum dependence of the four-gluon vertex



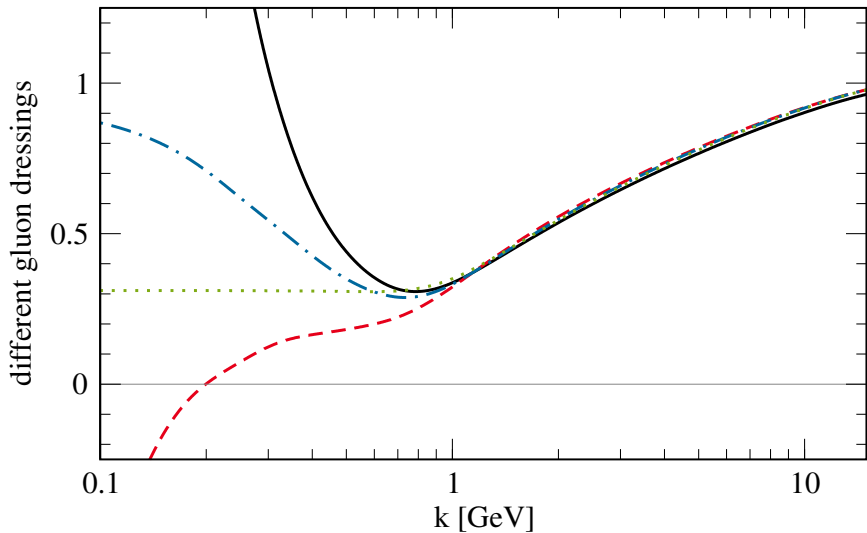
AKC, Fister, Mitter, Pawlowski, Strodthoff, 2016

# Four-gluon vertex at the symmetric point



AKC, Fister, Mitter, Pawłowski, Strodthoff, 2016

# Regulator dressing



AKC, Fister, Mitter, Pawłowski, Strodthoff, 2016